

# INVESTIGATION ON THE INFLUENCE OF GAS PRESSURE ON CO<sub>2</sub> ARC CHARACTERISTICS IN HIGH-VOLTAGE GAS CIRCUIT BREAKERS

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**Abstract.** CO<sub>2</sub> is identified as a promising alternative gas of SF<sub>6</sub>. The magnetohydrodynamics (MHD) arc model is established for a CO<sub>2</sub> circuit breaker. The influence of gas pressure is studied. The simulations are carried out for 0.5 MPa, 0.7 MPa and 0.9 MPa absolute filling pressure, allowing predictions of pressure and temperature distributions. The arc time constant  $\theta$  and the power loss coefficient  $Q$  is extracted. The thermal interruption capability is estimated to grow with increasing filling pressure.

**Keywords:** CO<sub>2</sub>, gas pressure, circuit breaker, Mayr equation.

## 1. Introduction

SF<sub>6</sub> is currently widely used in high-voltage gas circuit breakers because of its high dielectric strength and excellent arc interruption properties [1]. However, SF<sub>6</sub> is found to have extremely high global warming potential (approximately 23900 times that of CO<sub>2</sub>) [2]. Among the alternative gases of SF<sub>6</sub>, CO<sub>2</sub> and its mixtures are considered as the promising options and there has been some real products filled with pure CO<sub>2</sub> gas [3, 4].

Although the interruption performance of CO<sub>2</sub> is inferior to SF<sub>6</sub>, the gap might become smaller with optimization design of CO<sub>2</sub> circuit breakers based on improved understanding. When CO<sub>2</sub> gas is used as an arc-quenching medium in high-voltage switchgear, the gas pressure may be perhaps several atmospheric pressures. In this paper, the influence of the gas pressure on the interruption performance of the CO<sub>2</sub> circuit breakers is investigated.

## 2. Simulation model

A 252 kV puffer-type SF<sub>6</sub> circuit breaker with a short-circuit current interrupting capacity of 63 kA is employed in this paper. The dynamic arc behaviors are simulated by adopting a 2D axisymmetric MHD arc model, which takes account of the Joule heating, the electromagnetic effect (the Lorentz force), the radiation loss and the turbulence enhanced mass, momentum and energy transport. The MHD model is introduced in detail in our previous paper [5].

The plasma medium is assumed to satisfy the conditions of the local temperature equilibrium (LTE) and local chemical equilibrium (LCE). The nozzle ablation process is not considered in this paper. The thermodynamic and transport properties of CO<sub>2</sub> are taken from the study of Tanaka [6]. The net emission coefficient of CO<sub>2</sub> is taken from the study of Billoux [7]. The

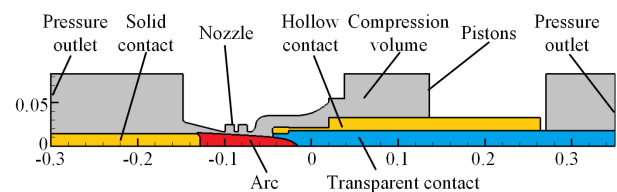


Figure 1. The simplified geometry and computational domain of the 252 kV puffer-type SF<sub>6</sub> circuit breakers

simplified geometry and computational domain of the circuit breakers are presented in Fig1.

The commercial computational software ANSYS FLUENT is used to model the circuit breaker and solve the governing equations. The boundary and initial conditions of the simulation model are listed as follows: (1) The circuit breakers are filled with pure CO<sub>2</sub> at an absolute pressure of 0.5 MPa, 0.7 MPa and 0.9 MPa; the initial temperature is set to 300 K; (2) The movement velocity of the moving contacts is piecewise linearized from the action characteristic experiment result; (3) The electric potential of the solid contact is set at zero; (4) The transparent contact hypothesis is adopted as in some papers; (5) The short-circuit current of 31.5 kA (50 Hz) rms and arcing time of 9.95 ms are adopted.

## 3. Temperature characteristics

The temperature distribution in the arc extinguish chamber is an important indicator of the interruption process. Fig. 2 shows the temperature distributions in the arc area at 0.5 ms before the current-zero with different filling pressures. The highest temperature is located at the exit of the compression volume. This is because that there is a stagnation zone at the exit of the compression volume where the cold gas from the compression volume gets separated into two parts. The temperature at the arc core increases with the in-

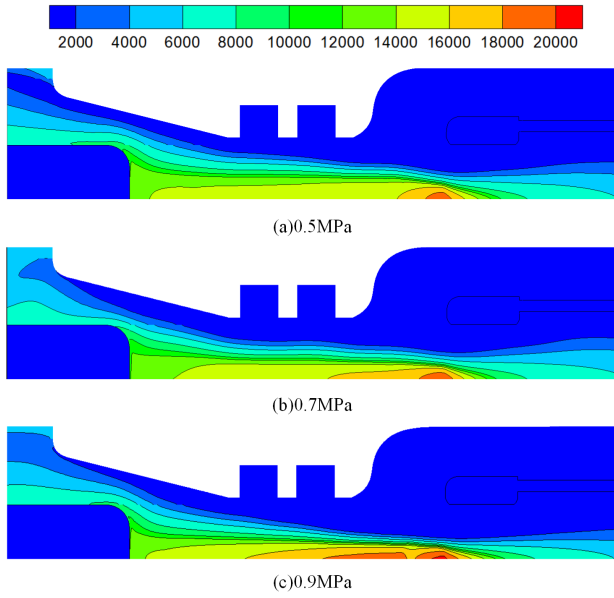


Figure 2. The temperature distributions in the arc area at 0.5 ms before the current-zero with different filling pressures. (Unit: K)

creasing filling pressure. The highest temperatures in the 0.5 MPa, 0.7 MPa and 0.9 MPa cases are 19.454 K, 19.803 K and 20.927 K.

Fig.3 describes the temperature distribution along the radial direction at the nozzle throat and at 0.5 ms before the current-zero with different filling pressures. It is clear that the temperature at the arc core becomes higher when the filling pressure is higher. However, the arc radius, defined as the position of the 4000 K isotherm in this paper as in some other papers, decreases with the increase of the filling pressure. What's more, there is an inflection point on the radial temperature profile around 8000 K in each of the three cases. Fig.4 gives the distribution of  $C_P$  along the radial direction at the nozzle throat and at 0.5 ms before the current-zero with different filling pressures. Coincidentally, the inflection points on the temperature profiles corresponds well with the three peaks on the  $C_P$  profiles for each cases. The peak value of  $C_P$  here is caused by the dissociation of CO and will have detrimental effects on arc's thermal recovery, because this effect will broaden the radiation reabsorption region and make the high temperature arc core less distinctive.

#### 4. Pressure characteristics

Fig.5 gives the pressure distributions in the arc area at 0.5 ms before the current-zero with different filling pressures. There is a distinctive stagnation zone at the exit of the compression volume in each case, with an obvious high pressure region. In the 0.5 MPa case, an arresting shock wave occurs between the solid contact and the nozzle divergence cone. However, in the 0.7 MPa and 0.9 MPa cases, the shock wave at this

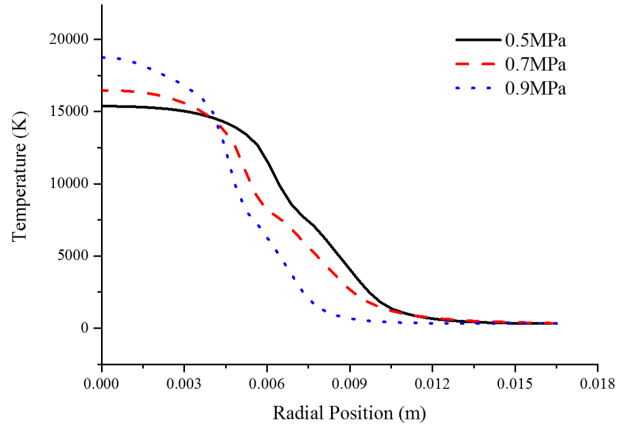


Figure 3. The temperature distribution along the radial direction at the nozzle throat ( $x=-0.07$  m) and at 0.5 ms before the current-zero with different filling pressures.

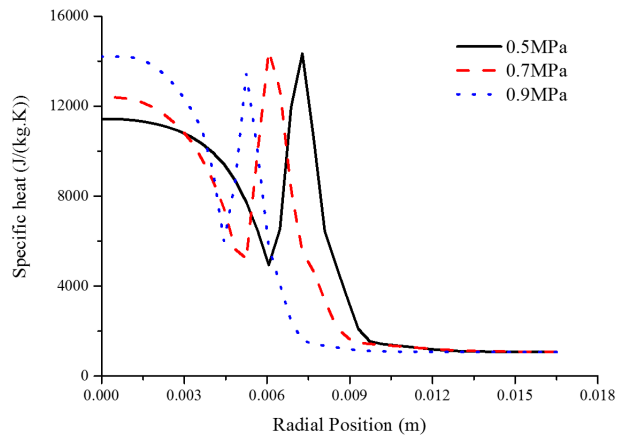


Figure 4. The distribution of the specific heat ( $C_P$ ) along the radial direction at the nozzle throat ( $x=-0.07$  m) and at 0.5 ms before the current-zero with different filling pressures.

location is much less stronger. Therefore, the filling pressure not only has an impact on the pressure values, but also affects the flow field during the interruption process. Besides, a slight shock wave also shows up inside the hollow contact in each case, which is caused by the contraction-expansion flow field structure.

#### 5. Current zero characteristics

The arc time constant  $\theta$  and the power loss coefficient  $Q$  are commonly used to evaluate the interruption capability of a circuit breaker, and generally give a believable result. Mayr equation is expressed as follow:

$$\frac{1}{g} \frac{dg}{dt} = \frac{1}{\theta} \left( \frac{u \cdot i}{Q} - 1 \right) \quad (1)$$

Then the time constant  $\theta$  and the power loss coefficient  $Q$  are expressed respectively:

$$\theta = \frac{g_1 i_2^2 - g_2 i_1^2}{i_1^2 \frac{dg_2}{dt} - i_2^2 \frac{dg_1}{dt}} \quad (2)$$

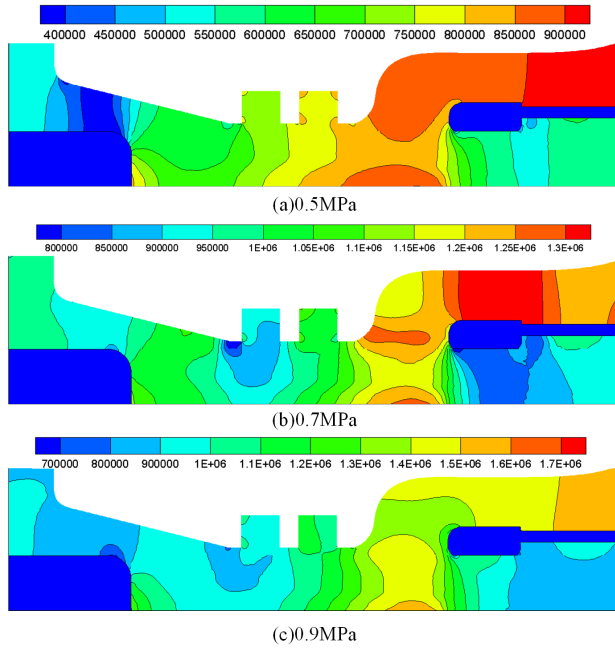


Figure 5. The pressure distributions in the arc area at 0.5 ms before the current-zero with different filling pressures. (Unit: Pa)

$$Q = \frac{i_2^2 \frac{dq_1}{dt} - i_1^2 \frac{dq_2}{dt}}{g_2 \frac{dq_1}{dt} - g_1 \frac{dq_2}{dt}} \quad (3)$$

In this paper, we assume that the time constant  $\theta$  and the power loss coefficient  $Q$  keep constant within a simulation time step. Thus, the values of  $\theta$  and  $Q$  at a specific time can be obtained by using two adjacent time steps as the two points required in equation (2) and (3). By this means, the variation of  $\theta$  and  $Q$  over time is available. The  $\theta$  and  $Q$  at the current zero point are obtained and shown in Fig.6. With the increase of the filling pressure, the time constant doesn't show much change with a value around 2.52  $\mu$ s. However, the power loss coefficient seems to increase monotonically with the increasing filling pressure.

In order to qualitatively evaluate the interruption capability of the circuit breaker, the post-arc current is predicted by using Mayr equation based on the calculated data of  $\theta$  and  $Q$  at the current zero point. The RRRV and relative  $di/dt$  are varied to find the interruption limit. The results for the 0.5 MPa case are shown in Fig.7 and Fig.8. In Fig.7, the relative  $di/dt$  is fixed as 1. In Fig. 8, the RRRV is set as 2 kV/ $\mu$ s. It is clearly seen that with increasing RRRV or relative  $di/dt$ , the post-arc current increases effectively, and when RRRV or relative  $di/dt$  is beyond the critical value, the current rises sharply and the arc re-ignition happens. The critical RRRV in the 0.5 MPa, 0.7 MPa and 0.9 MPa filling pressure cases is determined to be approximately 0.63 kV/ $\mu$ s, 0.74 kV/ $\mu$ s and 0.81 kV/ $\mu$ s respectively. The critical relative  $di/dt$  in the 0.5 MPa, 0.7 MPa and 0.9 MPa filling pressure cases is determined to be approximately 31

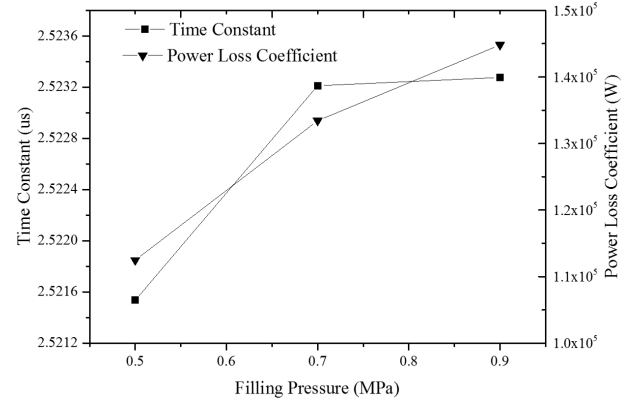


Figure 6. The time constant  $\theta$  and the power loss coefficient  $Q$  at the current zero point with different filling pressures

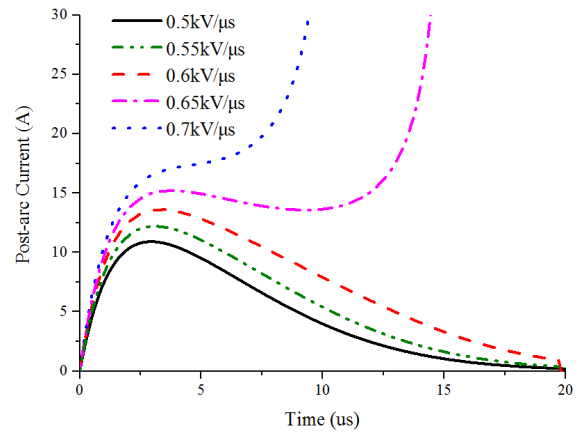


Figure 7. Variation of the post-arc current for different values of RRRV with the filling pressure of 0.7 MPa.

## 6. Conclusions

The performance of CO<sub>2</sub> gas as an interruption medium in high-voltage circuit breakers is investigated with different filling pressures. It is found that the highest temperature increases with the increasing filling pressure while the arc radius decreases with the increase of filling pressure. Besides, the arc temperature distribution has a close relationship with the specific heat ( $C_p$ ) of the filling gas. The filling pressure not only has an impact on the pressure values, but also affects the flow field during the interruption process. The post-arc current is predicted by using Mayr equation based on the calculated data of  $\theta$  and  $Q$  at the current zero point. For the 0.5 MPa, 0.7 MPa and 0.9 MPa filling pressure cases, the critical RRRV is determined to be 0.63 kV/ $\mu$ s, 0.74 kV/ $\mu$ s and 0.81 kV/ $\mu$ s and the critical relative  $di/dt$  is determined to be 31

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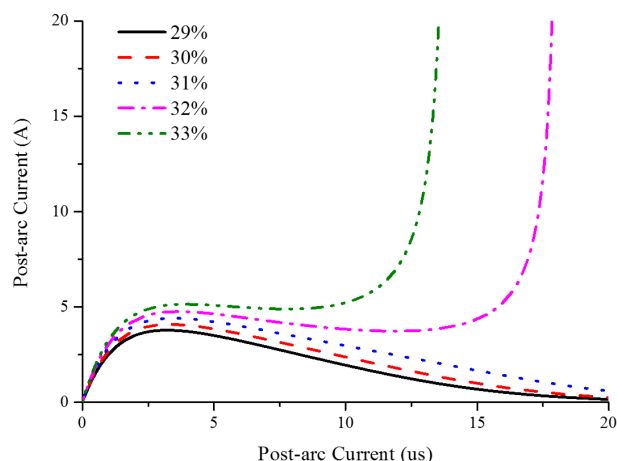


Figure 8. Variation of the post-arc current for different values of relative  $di/dt$  with the filling pressure of 0.5 MPa.

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