# EXPERIMENTAL AND SIMULATION STUDIES ON THE VOLTAGE DROP OF ARC IN LOW-VOLTAGE CIRCUIT BREAKER

J. Lu<sup>a</sup>, G. Déplaude<sup>b</sup>, P. Freton<sup>a</sup>, J-J. Gonzalez<sup>a,\*</sup>, P. Joyeux<sup>b</sup>

<sup>a</sup> Laboratory on Plasma and Conversion of Energy, University of Toulouse, CNRS, INPT, UPS., 118 Route de Narbonne, F-31062 Toulouse cedex 9, France

<sup>b</sup> Hager Electro SAS, 132 Bd d'Europe, 67215 Obernai, France

\* gonzalez@laplace.univ-tlse.fr

**Abstract.** In low voltage circuit breaker (LVCB) apparatus, a current limitation is performed by increasing the arc voltage. This increase is mainly realized in the splitters plates of the arc chamber by additional drop voltages due to anode and cathode sheaths regions. The consideration of the voltage drops near-electrodes regions is so one of the most important mechanism to improve the description of the arc behavior in LVCB. In this paper, the arc voltage evolution has been studied by experimental and simulation by considering a simple geometry constituted by two rails runner with one or two splitters plates. One magneto hydrodynamic model in three dimensions (3D) was developed to simulate the arc motion and the arc splitting process. In order to compare with the model results, experimental tests have been carried out.

Keywords: Low-voltage circuit breaker, voltage drop, splitters, experiment, simulation.

# 1. Introduction

The principle of LVCB consists to create a current limitation by increasing the arc voltage. The increase of the voltage is carried out by various mechanisms: increase of the losses, arc elongation and multiplication of the voltage drops in the splitters. To quantify and better describe these mechanisms, studies are conducted through models and experiments [1-4]. The most theoretical approach studied and applied in LVCB is the one proposed by Lindmayer [2]. This method allows considering an additional drop voltage through the resistivity in the layers. Another approach consists to describe the sheath with a nonequilibrium approach [5]. Nevertheless this method is difficult to implement for a real configuration with several splitters plates. Some authors [6-10] studied numerically the splitting process. Considering a nonlinear permeability of ferromagnetic material's they focus their studies on the influence of eddy currents [9] and on the influence of metal vapors coming from iron splitter plate erosion [10]. In those studies, the cathode voltage was treated by U-J curves and the anode arc root description was determined with the LTE-diffusion model of Lowke and Tanaka [11]. Another group has studied the arc phenomena in LVCB by an imaging system applied to a flexible test apparatus [12-15]. The studies were devoted to the influence of vent aperture size on re-ignition and arcing phenomena. They also proposed a modified method to consider the voltage due to the sheath regions. The ability to predict the arc extinction was analyzed and compared [16] with the Lindmayer approach [2]. The most complete study was proposed by Iturregi et al. [17]. They designed a 3D model in LVCB to analyze

the arc behavior. Nevertheless incoherencies appear in the comparison between one and two splitters plate: the additional drop voltage due to a second splitter is not present in the total voltage evolution versus time. The method is not clearly described as the authors used directly the @Ansys module which includes the Lindmayer approach. In this paper, the process of arc splitting is analyzed by theoretical and experimental works. Experimental tests have been realized using one and two splitters: high speed camera and electrical measurements. For the model we have implemented the Lindmayer approach and used the same geometries of Iturregi et al. [17].

# 2. Numerical method

#### 2.1. Hypotheses

The following assumptions have been used:

- (1) The air medium plasma is assumed to be in Local Thermodynamic Equilibrium (LTE).
- (2) The plasma is a Newtonian fluid and the flow is laminar.
- (3) Vapors from the walls and electrodes are not taken into account.
- (4) Radiation is treated using the net emission coefficient method.

#### 2.2. Equations

Based on the Fluent software, we have implemented a fluid model to simulate electrical arc behavior in a LVCB. It can be characterized by the macroscopic fluid quantities: velocity, pressure, enthalpy. All these quantities are obtained from the Navier-Stokes equations with Lorentz forces and homic heating and completed with Maxwell equations to consider electromagnetism effects. The magnetic field is obtained from the current path in the runners and plasma through the vector potential resolution and a BiotSavart formulation is used at the boundaries. In order to represent the additional voltage drop for the runners and splitters, an effective electrical conductivity is defined [2] based on experimental works [6, 8]. The voltage evolution versus the current density is plotted in the Figure 1. These characteristics (a-b-c-d) suggested by Lindmayer et al. [2] in air medium with copper runners enable some tuning of the model to adjust the voltage, they represent the current ability to jump on the splitter plate.



Figure 1. Evolution of voltage drop versus current density [3]

#### 2.3. Boundary conditions

"Classical" boundary conditions are used [1]. The magnetic field is calculated from the vector potentials but the BiotSavart formulation is used to close the system resolution.

### 2.4. Geometries

The geometries for the simulation are presented in the Figure 2. The dimensions of the chambers are  $40 \times 2.5 \times 11 \text{ mm} (xyz)$  and the rails are  $40 \times 2.5 \times$ 1.5 mm (xyz). It should be noted that the dimension of the splitters plate are not the same: for one splitter:  $20 \times 2.5 \times 2 \text{ mm} (xyz)$  and for two splitters:  $20 \times$  $2.5 \times 1 \text{ mm} (xyz)$ . D1 = 8 mm represents the distance between the two runners; D2 = 3 mm is the distance between runners and splitters, for the third geometry D3 = 2 mm. These dimensions are far from real LVCB but chosen to correspond with Iturregi et al. [17]. A uniform grid is used in the geometry leading to 1.12 million cells.

### 3. Simulation results

We present the main results obtained with the 3D model applying the Lindmayer approach. The arc



Figure 2. Geometries used for the simulation (From left to right: 0 splitter, 1 splitter and 2 splitters).

splitting behavior is shown through the current densities in Figure 4 for one and two splitter plates. Figure 3 presents the voltage evolution for the two cases. The calculation time step is 10  $\mu$ s and the input current is 50 A DC. The curve "c" of Figure 1 was chosen to simulate the sheath contribution. There is no symmetry plane in our simulation. In order to reduce the calculation time, the arc is ignited 1.5 mm front of splitters plates.



Figure 3. Voltage drop in the simulation with 1 and 2 splitters.

In case of one splitter: At  $t = 50 \,\mu\text{s}$  the arc starts to bend and moves toward the splitters. Due to the arc elongation the voltage increases progressively from 61 V at  $t = 50 \,\mu\text{s}$  to its maximum value 94 V at  $t = 0.19 \,\mathrm{ms}$ . Then the voltage decreases. At  $t = 0.19 \,\mathrm{ms}$  the drop voltage is due to the effects of the arc elongation and to the additional voltage sheathes. After time  $t = 0.19 \,\mathrm{ms}$ , the current flows progressively into the splitter leading to a diminution of the current path and of the total voltage. The current path goes into the splitter totally at t = 0.21 ms.At t = 0.25 ms there is only one path of current and the voltage drop decreases to 81 V. The contribution of the splitter sheath on the total voltage is around 30 V in the following case. In case of two splitters, the presented times are not the same; indeed the

Boundary conditions	Momentum	Enthalpy	Scalar Potential	Vector Potential
Cathode	$v=0 \mathrm{m/s}$	Heat Transfer	Continuity	BiotSavart
Anode	$v=0 \mathrm{m/s}$	Heat Transfer	Continuity	BiotSavart
In	/	dT/dn=0	I=50 A	/
Out	$v=0 \mathrm{m/s}$	dT/dn=0	$0 \mathrm{V}$	/
Vent In	$v=0 \mathrm{m/s}$	$300\mathrm{K}$	dV/Dn=0	BiotSavart
Vent Out	P=1 atm	Convection	dV/Dn=0	BiotSavart
Walls	$v=0 \mathrm{m/s}$	$300\mathrm{K}$	dV/Dn=0	/
Splitters	v=0m/s	Heat Transfer	Continuity	BiotSavart

Table 1. Boundary conditions used in the simulation model.



Figure 4. Simulated current densities using curve c for one (left) and two (right) splitters.

arc motion, the splitting process time and behavior are different. At t = 0.14 ms, the arc column bends around the splitter plates, the arc voltage increases up to 76.6 V. Compared to one splitter case, we can observe that at t = 0.15 ms the arc begins to jump on the splitters. From this time, two arc roots appear on the splitters plates, the current density through the splitter plates increases and the current path out of the splitters plates gradually fades. Compared with the case of one splitter, the arc voltage evolution is faster. At time  $t = 0.25 \,\mathrm{ms}$  the difference on the total arc voltage between the two cases is only 4 V. The comparison is difficult to make as the current paths and arc lengths are not the same. According to the results with one splitter an additional voltage U = 30 Vbetween the two cases should exist. Nevertheless the arc positions are different, the medium temperature is different and the arc behavior not the same as shown by the voltage evolutions Figure 3.

## $\begin{array}{ccc} \text{mp on} & 50 \times 15 \end{array}$

following:  $126 \times 15 \times 20 \,\mathrm{mm} \,(xyz)$  for the chamber,  $126 \times 15 \times 1 \,\mathrm{mm}$  (xyz) for the runners and  $50 \times 15 \times 1 \,\mathrm{mm}$  for the splitter plate. There is a wire located in front of the splitter to ignite the electrical arc. The experimental setup is composed by a generator which can product a prospective peak current up to 13 kA with a maximum charging voltage of 600 V in AC, a high-speed camera (Photron SA5), a differential voltage probe and a Rogowski coil for the electrical measurements. Figure 5 presents the current evolution and the total voltages versus time for zero and one splitter. The first peak of voltage at  $t = 0.25 \,\mathrm{ms}$  characterizes the arc ignition. After that, the arc moves toward to the upper side of the chamber. During the arc motion the voltage is nearly constant 60 V until t = 2 ms. At this time the arc arrives in front of the splitter plate. In the case with

The dimensions of the experimental setup are the

4. Experimental results



Figure 5. Experimental voltage drop and total current. Cases with one and without splitter in the chamber.



Figure 6. Example of arc behavior in front of one and two splitters plates obtained by high-speed camera.

one splitter, arc voltage increases by squeezing against the splitter, by touching the lower edge of the splitter until arc was totally on the splitter from t = 2 ms to t = 3 ms. The voltage drop contribution dues to the splitter can be calculated by the difference between the case without splitter. The value is nearly 25 V. This sheath contribution is in the same order of magnitude than the theoretical case, nevertheless we have seen between the two theoretical cases (one and two splitters) that the comparison is difficult due to the tortuous current paths or arc behavior in the plasma.

Finally we present Figure 6 two images obtained by the high speed camera to illustrate the arc behavior observed experimentally. In the upper picture the arc is bent in front of the splitter as in the down picture the arc is segmented by the presence of the two splitters plates.

## 5. Conclusion

A 3D model was developed and the effective electrical model of conductivity suggested by Lindmayer [2] introduced in the developments to characterize the additional voltage due to the splitter plate sheathes. The feasibility is demonstrated in cases of one and two splitters plates. Interesting results are obtained on the voltage evolutions. In the same time investigations are made to characterize the arc behavior in an experimental setup by high speed camera and electrical measurements. The theoretical and experimental configurations are at this time different nevertheless the same order of magnitude was observed with an additional drop voltage around 20–30 V for each splitter plate. The next step of our study will consist to use

Voltage drops in LVCB

the same geometry for the model and the experience and to perform one parametric study on the number of splitters plates.

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